

TABLE 1. Thermoelastic properties of brain matter [Lin, 1977a].

Specific Heat, $c_h$	0.88 cal/g-°C (0.21 J/g-°C)
Density, $\rho$	1.05 g/cm <sup>3</sup>
Coefficient of Thermal Expansion, $\alpha$	$4.1 \times 10^{-5}/^\circ\text{C}$
Lame's Constant, $\lambda$	$2.24 \times 10^{10}$ dyn/cm <sup>2</sup>
Lame's Constant, $\mu$	$10.52 \times 10^3$ dyn/cm <sup>2</sup>
Bulk Velocity of Propagation, $c_1$	$1.46 \times 10^5$ cm/s

If we choose the median of these two curves, it can be seen that, for a guinea pig ( $a = 2$  cm), the predicted fundamental frequency is 45 kHz, which is close to that measured by Chou *et al.* [1975]. For a typical cat ( $a = 3$  cm) the measured cochlear microphonic frequency is 38 kHz. Figure 1 shows a computed frequency of 30 kHz. The computed frequency for an infant ( $a = 5$  cm) and an adult human being ( $a = 7$  cm) are 18 and 13 kHz, respectively. Although experimental data that are directly applicable to this case are not available, related results have indicated that a necessary condition for auditory perception by adults seems to be the ability to hear sounds above 5 to 8 kHz [Frey, 1961; Rissmann and Cain, 1975].

3.2. Peak pressure and displacement.

Extensive numerical computations have shown that, although the resultant waveforms are qualitatively similar, both pressures and displacements that are computed on the basis of the constrained-surface formulation are consistently higher than those calculated using the stress-free expressions. Moreover, better agreement between theory and experiment exists for the constrained-boundary formulation. The following presentation is therefore specifically related to the constrained-boundary expressions given by equations (13) to (20).

Figures 2 and 3 show the peak pressure and displacement at  $r = 0, 1$  and 2 cm for a sphere of 2-cm radius that stimulates the head of a guinea pig under 2450-MHz radiation. In this case the approximate microwave absorption pattern is obtained by setting  $N = 3$ . These waveforms are evaluated for  $t_0 = 10 \mu\text{sec}$ . It is seen that the pressure is the highest at the center of the head and has a maximum value of 4.08 dyne/cm<sup>2</sup> for a peak absorption rate of 1 W/g (which corresponds to a power density 445 mW/cm<sup>2</sup> of incident energy [Johnson and Guy, 1972]) at the end of the pulse and then oscillates around a constant average value in the absence of elastic damping. As expected, the displacement is zero both at the center and at the surface of the sphere for the constrained-boundary condition. The peak displacement at  $r = 1.0$  cm is about  $2.16 \times 10^{-13}$  meters. Note that a high-frequency component is superimposed on top of the fundamental frequency, which correlates perfectly with that shown in Figure 1.

The peak pressure and displacement at  $r = 0, 1.5$  and 3.0 cm for a 3-cm sphere, which simulates the head of a cat exposed to 2450-MHz radiation, are illustrated in Figures 4 and 5. In this case the approximate microwave absorption pattern is obtained by setting  $N = 6$ . Again, the pressure is the highest at the center, and has a value of 3.69 dyne/cm<sup>2</sup> for a peak absorption rate of 1 W/g or 589 mW/cm<sup>2</sup> of incident energy. Quantitative comparisons with two series of experimental threshold determinations [Guy *et al.* 1975; Rissmann and Cain, 1975] indicate that the computed threshold is extremely close to the measured ones [Lin, 1977b]. The displacement in the model of the cat's head is about  $1.51 \times 10^{-13}$  meters. The fundamental frequency of oscillation ob-

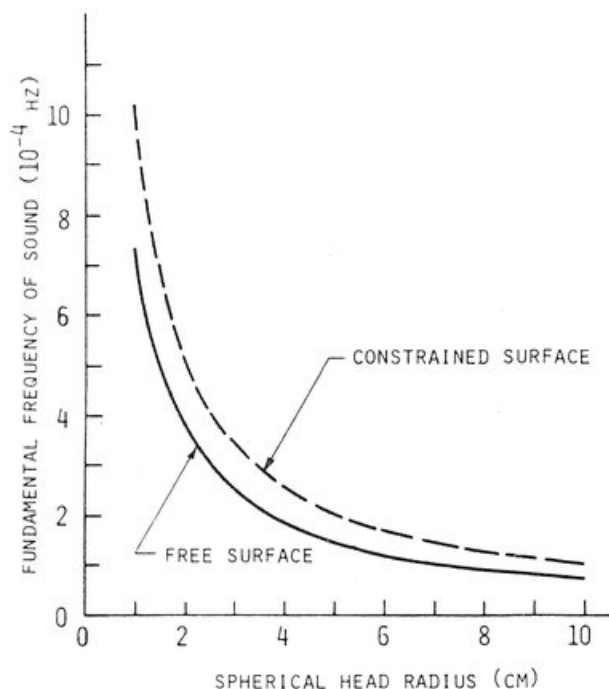


Fig. 1. Predicted fundamental frequency of vibration of a spherical head in a pulse-moderated microwave field.